



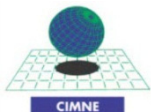
# Continuum Mechanics

## Chapter 5 Stresses

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# Forces

## Chapter 5 · Stresses

1. Forces
2. Cauchy's stress theorems
3. Stress tensors

# Forces

## Forces

We consider two types of forces that may act on a continuum body: **body** (or **mass** or **volume** or **internal**) **forces** and **surface forces**.

- **Body** (or **mass** or **volume** or **internal**) forces: Forces acting in the *volume* of a continuum medium. Typical examples of body forces are the gravity forces or the electromagnetic forces.
- **Surface forces**: Forces acting on the *surface* of a continuum medium due to the interaction with other bodies or the environment. Typical examples of surface forces are the contact forces or applied loads.

# Body Forces

## Body Forces

**Body** (or **mass** or **volume** or **internal**) **forces** may be characterized by the **body forces per unit of mass** vector, denoted as  **$\mathbf{b}$** .

*Spatial* and *material descriptions* of the **body forces per unit of mass** vector, denoted as  **$\mathbf{b}(\mathbf{x}, t)$**  and  **$\mathbf{B}(\mathbf{X}, t)$** , respectively, take the form,

$$\mathbf{b} = \mathbf{b}(\mathbf{x}, t) = \mathbf{b}(\boldsymbol{\varphi}(\mathbf{X}, t), t) = \mathbf{B}(\mathbf{X}, t)$$

$$\mathbf{b} = \mathbf{B}(\mathbf{X}, t) = \mathbf{B}(\boldsymbol{\varphi}^{-1}(\mathbf{x}, t), t) = \mathbf{b}(\mathbf{x}, t)$$

# Body Forces

## Differential Body Forces

The **differential body** (or **mass** or **volume** or **internal**) force acting in a differential of volume  $dv$  in the *spatial configuration* takes the form,

$$d\mathbf{f}_v = \rho(\mathbf{x}, t) \mathbf{b}(\mathbf{x}, t) dv$$

The **differential body** (or **mass** or **volume** or **internal**) force acting in a differential of volume  $dV$  in the *material configuration* takes the form,

$$d\mathbf{f}_v = \rho_0(\mathbf{X}) \mathbf{B}(\mathbf{X}, t) dV$$

# Body Forces

## Total Body Forces

The **total body** (or **mass** or **volume** or **internal**) **forces** acting in a *spatial volume*  $v$  of a continuum body, at a time  $t$ , may be written as,

$$\mathbf{F}_v = \int_v \rho(\mathbf{x}, t) \mathbf{b}(\mathbf{x}, t) dv$$

The **total body** (or **mass** or **volume** or **internal**) **forces** acting in a *material volume*  $V$  of a continuum body, at a time  $t$ , may be written as,

$$\mathbf{F}_v = \int_V \rho_0(\mathbf{X}) \mathbf{B}(\mathbf{X}, t) dV$$

# Surface Forces

## Differential Surface Forces

The **differential surface force** acting on a differential of area  $ds$  on the *spatial configuration* takes the form,

$$d\mathbf{f}_s = \mathbf{t}(\mathbf{x}, t) ds$$

where the **Cauchy (or true) traction** vector, denoted as  $\mathbf{t}(\mathbf{x}, t)$ , represents the *spatial description* of the surface force per unit of *spatial surface*.

The **differential surface force** acting on a differential of area  $dS$  on the *material configuration* takes the form,

$$d\mathbf{f}_s = \mathbf{T}(\mathbf{X}, t) dS$$

where the **first Piola-Kirchhoff (or nominal) traction** vector, denoted as  $\mathbf{T}(\mathbf{X}, t)$ , represents the *material description* of the surface force per unit of *material surface*.

# Surface Forces

## Total Surface Forces

The **total surface forces** acting on a *spatial surface*  $\partial v$  of a continuum body, at a time  $t$ , may be written in terms of the **Cauchy (or true) traction** vector as,

$$\mathbf{F}_{\partial v} = \int_{\partial v} \mathbf{t}(\mathbf{x}, t) ds$$

The **total surface forces** acting on a *material surface*  $\partial V$  of a continuum body, at a time  $t$ , may be written in terms of the **first Piola-Kirchhoff (or nominal) traction** vector as,

$$\mathbf{F}_{\partial V} = \int_{\partial V} \mathbf{T}(\mathbf{X}, t) dS$$



# Forces

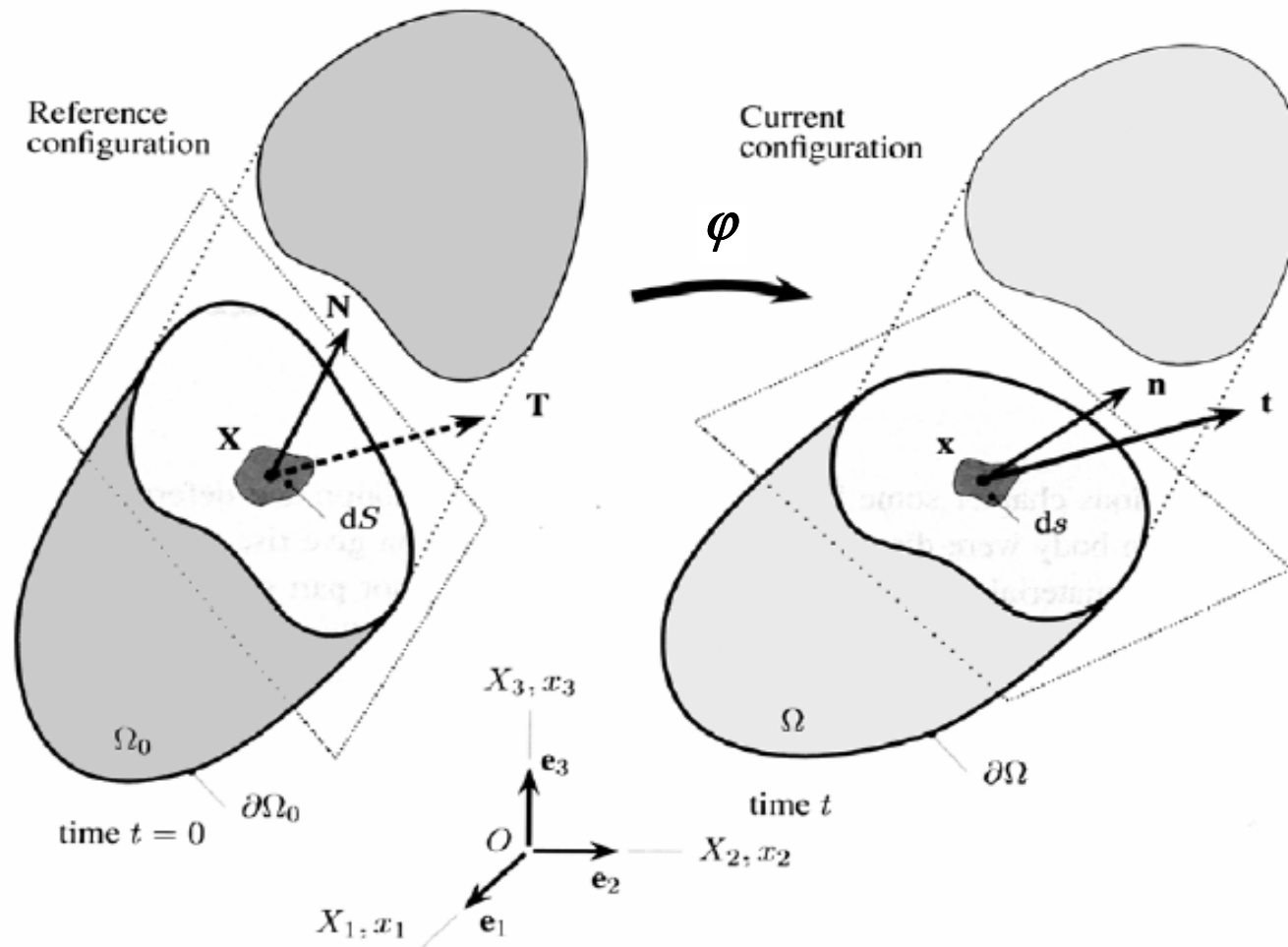
## Body Forces

$$\mathbf{F}_v = \int_v \rho(\mathbf{x}, t) \mathbf{b}(\mathbf{x}, t) dv = \int_V \rho_0(\mathbf{X}) \mathbf{B}(\mathbf{X}, t) dV$$

## Surface Forces

$$\mathbf{F}_{\partial v} = \int_{\partial v} \mathbf{t}(\mathbf{x}, t) ds = \int_{\partial V} \mathbf{T}(\mathbf{X}, t) dS$$

# The Traction Vector Picture



# Cauchy Stress Theorems

## First Cauchy Stress Theorem

The **Cauchy (or true) traction** vector at a spatial point  $\mathbf{x}$ , at a given time  $t$ , on a *spatial surface* with *unit outward normal*  $\mathbf{n}$  at the spatial point  $\mathbf{x}$ , is only a function of the *spatial point*, the *time*  $t$  and the *unit outward normal* at the spatial point  $\mathbf{x}$  at the time  $t$ ,

$$\mathbf{t} = \mathbf{t}(\mathbf{x}, t, \mathbf{n}), \quad t_a = t_a(\mathbf{x}, t, \mathbf{n})$$

# Cauchy Stress Theorems

## First Cauchy Stress Theorem

The **first Piola-Kirchhoff** (or **nominal**) **traction** vector at a material point  $\mathbf{X}$ , at a given time  $t$ , on a *material surface* with *unit outward normal*  $\mathbf{N}$  at the material point  $\mathbf{X}$ , is only a function of the *material point*, the *time*  $t$  and the *unit outward normal* at the material point  $\mathbf{X}$  at the time  $t$ ,

$$\mathbf{T} = \mathbf{T}(\mathbf{X}, t, \mathbf{N}), \quad T_a = T_a(\mathbf{X}, t, \mathbf{N})$$

# Cauchy Stress Theorems

## Second Cauchy Stress Theorem

The **Cauchy** (or **true**) **traction** vector at a spatial point  $\mathbf{x}$ , at a given time  $t$ , on a *spatial surface* with *unit outward normal*  $\mathbf{n}$  at the spatial point  $\mathbf{x}$ , is a *linear* function of the *unit outward normal* at the spatial point  $\mathbf{x}$  at the time  $t$ , satisfying the so called **action-reaction principle**,

$$\mathbf{t} = \mathbf{t}(\mathbf{x}, t, \mathbf{n}) = -\mathbf{t}(\mathbf{x}, t, -\mathbf{n}), \quad t_a = t_a(\mathbf{x}, t, \mathbf{n}) = -t_a(\mathbf{x}, t, -\mathbf{n})$$

# Cauchy Stress Theorems

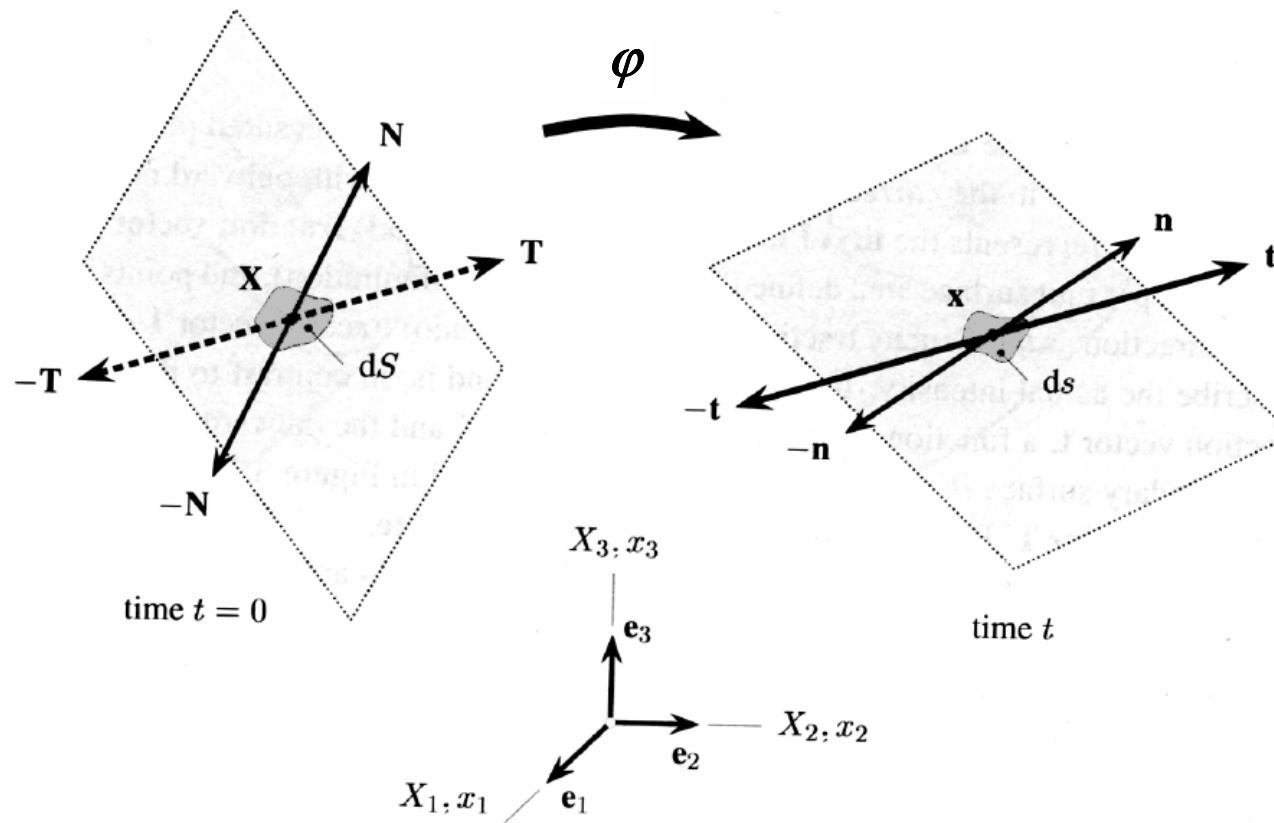
## Second Cauchy Stress Theorem

The **first Piola-Kirchhoff** (or **nominal**) **traction** vector at a material point  $\mathbf{X}$ , at a given time  $t$ , on a *material surface* with *unit outward normal*  $\mathbf{N}$  at the material point  $\mathbf{X}$ , is a *linear* function of the *unit outward normal* at the material point  $\mathbf{X}$  at the time  $t$ , satisfying the so called **action-reaction principle**,

$$\mathbf{T} = \mathbf{T}(\mathbf{X}, t, \mathbf{N}) = -\mathbf{T}(\mathbf{X}, t, -\mathbf{N}), \quad T_a = T_a(\mathbf{X}, t, \mathbf{N}) = -T_a(\mathbf{X}, t, -\mathbf{N})$$

# Cauchy Stress Theorems

## Second Cauchy Stress Theorem



# Stress Tensors

## Cauchy Stress Tensor

The **Cauchy** (or **true**) **stress** tensor, denoted as  $\boldsymbol{\sigma}$ , is a *symmetric* spatial second-order tensor, such that,

$$\mathbf{t}(\mathbf{x}, t, \mathbf{n}) = \boldsymbol{\sigma}(\mathbf{x}, t) \mathbf{n}, \quad t_a(\mathbf{x}, t, \mathbf{n}) = \sigma_{ab}(\mathbf{x}, t) n_b$$

## First Piola-Kirchhoff Stress Tensor

The **first Piola-Kirchhoff** (or **nominal**) **stress** tensor, denoted as  $\mathbf{P}$ , is a *non-symmetric* two-point second-order tensor, such that,

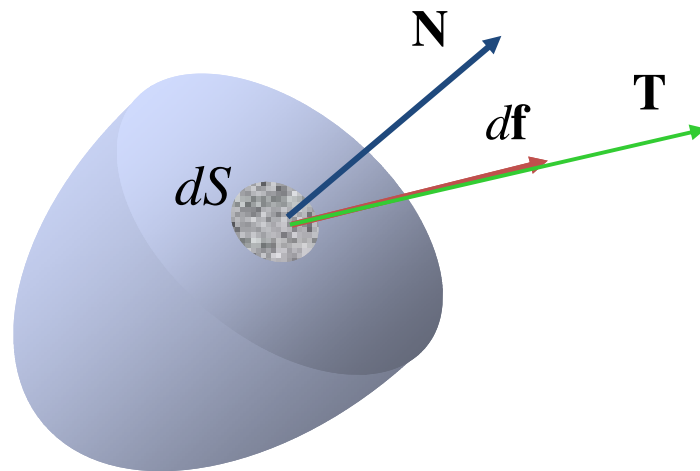
$$\mathbf{T}(\mathbf{X}, t, \mathbf{N}) = \mathbf{P}(\mathbf{X}, t) \mathbf{N}, \quad T_a(\mathbf{X}, t, \mathbf{N}) = P_{aA}(\mathbf{X}, t) N_A$$



# Stress Tensors

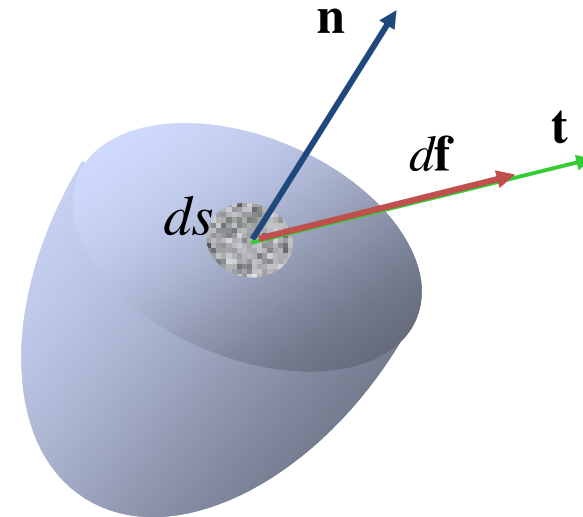
Reference Configuration

time  $t = 0$



Current Configuration

time  $t$



$$d\mathbf{f} = \mathbf{T}dS = \mathbf{t} ds$$

$$\mathbf{T} = d\mathbf{f}/dS$$

$$\mathbf{T} = \mathbf{T}(\mathbf{X}, t, \mathbf{N}) = \mathbf{P}(\mathbf{X}, t) \mathbf{N}$$

$$\mathbf{t} = d\mathbf{f}/ds$$

$$\mathbf{t} = \mathbf{t}(\mathbf{x}, t, \mathbf{n}) = \boldsymbol{\sigma}(\mathbf{x}, t) \mathbf{n}$$

# Piola Transformation

## Piola Transformation

The *Cauchy* (or *true*) *traction* vector and the *first Piola-Kirchhoff* (or *nominal*) *traction* vector are related through the expression,

$$d\mathbf{f} = \mathbf{t} ds = \mathbf{T} dS$$

Introducing the *Cauchy* (or *true*) *stress* tensor and the *first Piola-Kirchhoff* (or *nominal*) *stress* tensor, yields,

$$d\mathbf{f} = \boldsymbol{\sigma} \mathbf{n} ds = \mathbf{P} \mathbf{N} dS$$

and using *Nanson's* formula, given by,

$$\mathbf{n} ds = J \mathbf{F}^{-T} \mathbf{N} dS$$

yields the so called **Piola transformation**, given by,

$$\mathbf{P} = J \boldsymbol{\sigma} \mathbf{F}^{-T}, \quad P_{aA} = J \sigma_{ab} F_{bA}^{-T}$$

# Piola Transformation

## Piola Identity

Using the diverge theorem, the following useful identity holds,

$$\int_{\partial\Omega} \mathbf{n} ds = \int_{\partial\Omega} \mathbf{1} n ds = \int_{\Omega} \operatorname{div} \mathbf{1} dv = \mathbf{0}$$

Using Nanson's formula and the divergence theorem yields,

$$\int_{\partial\Omega} \mathbf{n} ds = \int_{\partial\Omega_0} \mathcal{J} \mathbf{F}^{-T} \mathbf{N} dS = \int_{\Omega_0} \operatorname{DIV} \left( \mathcal{J} \mathbf{F}^{-T} \right) dV = \mathbf{0}$$

And in local form, we obtain the so called **Piola identity** given by,

$$\operatorname{DIV} \left( \mathcal{J} \mathbf{F}^{-T} \right) = \mathbf{0}, \quad \left( \mathcal{J} F_{aA}^{-T} \right)_{,A} = 0$$

# Piola Transformation

## Piola Transformation

Using the **Piola transformation** and the **Piola identity** given by,

$$\mathbf{P} = J \boldsymbol{\sigma} \mathbf{F}^{-T}, \quad \text{DIV} \left( J \mathbf{F}^{-T} \right) = \mathbf{0},$$

the following expression holds,

$$\text{DIV} \mathbf{P} = \text{DIV} \left( J \boldsymbol{\sigma} \mathbf{F}^{-T} \right) = J \mathbf{F}^{-T} \text{DIV} \boldsymbol{\sigma} + \cancel{\boldsymbol{\sigma} \text{DIV} \left( J \mathbf{F}^{-T} \right)} = J \text{div} \boldsymbol{\sigma}$$

yielding the useful expression,

$$\text{DIV} \mathbf{P} = J \text{div} \boldsymbol{\sigma}, \quad P_{aA,A} = J \sigma_{ab,b}$$

# Piola Transformation

## Symmetry Restriction

The *Piola transformation* yields to the following relations between the *Cauchy* (or *true*) stress tensor and the *first Piola-Kirchhoff* (or *nominal*) stress tensor,

$$\mathbf{P} = J \boldsymbol{\sigma} \mathbf{F}^{-T}, \quad \boldsymbol{\sigma} = J^{-1} \mathbf{P} \mathbf{F}^T$$

The *symmetry* of the *Cauchy stress tensor*, i.e.,

$$\boldsymbol{\sigma} = \boldsymbol{\sigma}^T, \quad \sigma_{ab} = \sigma_{ab}^T = \sigma_{ba}$$

yields the following *symmetry restriction* on the *first Piola-Kirchhoff stress tensor*,

$$\mathbf{P} \mathbf{F}^T = \mathbf{F} \mathbf{P}^T, \quad P_{aA} F_{Ab}^T = P_{aA} F_{bA} = F_{aA} P_{Ab}^T = F_{aA} P_{bA}$$

# Stress Tensors

## Cauchy Stress Tensor

$$\boldsymbol{\sigma} = \boldsymbol{\sigma}^T, \quad \sigma_{ab} = \sigma_{ab}^T = \sigma_{ba}$$

$$\boldsymbol{\sigma} = J^{-1} \mathbf{P} \mathbf{F}^T, \quad \sigma_{ab} = J^{-1} P_{aA} F_{Ab}^T = J^{-1} P_{aA} F_{bA}$$

$$\operatorname{div} \boldsymbol{\sigma} = J^{-1} \operatorname{DIV} \mathbf{P}, \quad \sigma_{ab,b} = J^{-1} P_{aA,A}$$

## First Piola-Kirchhoff Stress Tensor

$$\mathbf{P} \mathbf{F}^T = \mathbf{F} \mathbf{P}^T, \quad P_{aA} F_{Ab}^T = P_{aA} F_{bA} = F_{aA} P_{Ab}^T = F_{aA} P_{bA}$$

$$\mathbf{P} = J \boldsymbol{\sigma} \mathbf{F}^{-T}, \quad P_{aA} = J \sigma_{ab} F_{bA}^{-T} = J \sigma_{ab} F_{Ab}^{-1}$$

$$\operatorname{DIV} \mathbf{P} = J \operatorname{div} \boldsymbol{\sigma}, \quad P_{aA,A} = J \sigma_{ab,b}$$

# Stress Tensor Components

## Cauchy Stress Tensor Components

The *Cauchy* (or *true*) *traction* vector at a *spatial point*, along the *Cartesian planes*, i.e. on planes with *unit vectors* along the Cartesian axes at the spatial configuration, read,

$$\mathbf{t}_1 = \boldsymbol{\sigma} \mathbf{e}_1 = \sigma_{11} \mathbf{e}_1 + \sigma_{21} \mathbf{e}_2 + \sigma_{31} \mathbf{e}_3$$

$$\mathbf{t}_2 = \boldsymbol{\sigma} \mathbf{e}_2 = \sigma_{12} \mathbf{e}_1 + \sigma_{22} \mathbf{e}_2 + \sigma_{32} \mathbf{e}_3$$

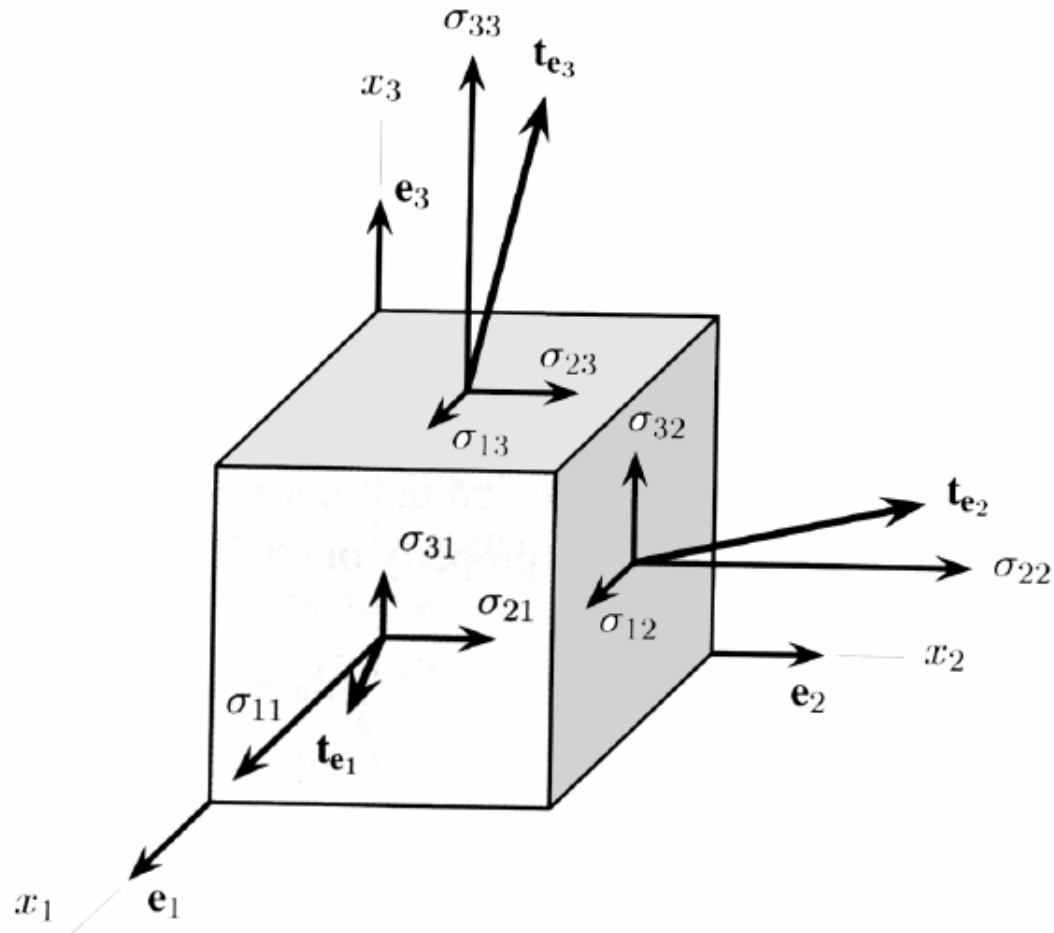
$$\mathbf{t}_3 = \boldsymbol{\sigma} \mathbf{e}_3 = \sigma_{13} \mathbf{e}_1 + \sigma_{23} \mathbf{e}_2 + \sigma_{33} \mathbf{e}_3$$

Note that the *ab*-component of the *Cauchy stress* tensor may be computed as,

$$\sigma_{ab} = \mathbf{e}_a \cdot \mathbf{t}_b = \mathbf{e}_a \cdot \boldsymbol{\sigma} \mathbf{e}_b$$

# Stress Tensor Components

## Cauchy Stress Tensor Components





# Stress Tensor Components

## Cauchy Stress Tensor Components

Using *index notation*, the *matrix of Cartesian components* of the *symmetric Cauchy (or true) stress tensor* takes the form,

$$[\boldsymbol{\sigma}] = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{12} & \sigma_{22} & \sigma_{23} \\ \sigma_{13} & \sigma_{23} & \sigma_{33} \end{bmatrix}$$

Being a *symmetric* second-order tensor, we may collect the six components into a vector of components, such that,

$$[\boldsymbol{\sigma}] = [\sigma_{11} \quad \sigma_{22} \quad \sigma_{33} \quad \sigma_{12} \quad \sigma_{13} \quad \sigma_{23}]^T$$

# Stress Tensor Components

## Cauchy Stress Tensor Components

Using *engineering notation*, the *matrix of Cartesian components* of the *symmetric Cauchy (or true) stress tensor* takes the form,

$$[\boldsymbol{\sigma}] = \begin{bmatrix} \sigma_x & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \sigma_y & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \sigma_z \end{bmatrix} = \begin{bmatrix} \sigma_x & \tau_{xy} & \tau_{xz} \\ \tau_{xy} & \sigma_y & \tau_{yz} \\ \tau_{xz} & \tau_{yz} & \sigma_z \end{bmatrix}$$

Being a *symmetric* second-order tensor, we may collect the six components into a vector of components, such that,

$$[\boldsymbol{\sigma}] = \left[ \sigma_x \quad \sigma_y \quad \sigma_z \quad \tau_{xy} \quad \tau_{xz} \quad \tau_{yz} \right]^T$$

## Other Stress Tensors

### Kirchhoff Stress Tensor

The **Kirchhoff stress** tensor, denoted as  $\boldsymbol{\tau}$ , is a *symmetric spatial* second-order tensor and it may be defined in terms of the *Cauchy stress* tensor as,

$$\boldsymbol{\tau} = J\boldsymbol{\sigma}, \quad \tau_{ab} = J\sigma_{ab}$$

### Second Piola-Kirchhoff Stress Tensor

The **second Piola-Kirchhoff stress** tensor, denoted as  $\mathbf{S}$ , is a *symmetric material* second-order tensor and it may be defined in terms of the *Kirchhoff stress* tensor as,

$$\mathbf{S} = \mathbf{F}^{-1}\boldsymbol{\tau}\mathbf{F}^{-T}, \quad S_{AB} = F_{Aa}^{-1}\tau_{ab}F_{bB}^{-1} = F_{Aa}^{-1}\tau_{ab}F_{Bb}^{-1}$$

# Push-forward / Pull-back operations

## Push-forward / Pull-back Operations

The *Kirchhoff stress* tensor may be viewed as the **push-forward** of the (contravariant) *second Piola-Kirchhoff stress* tensor, satisfying,

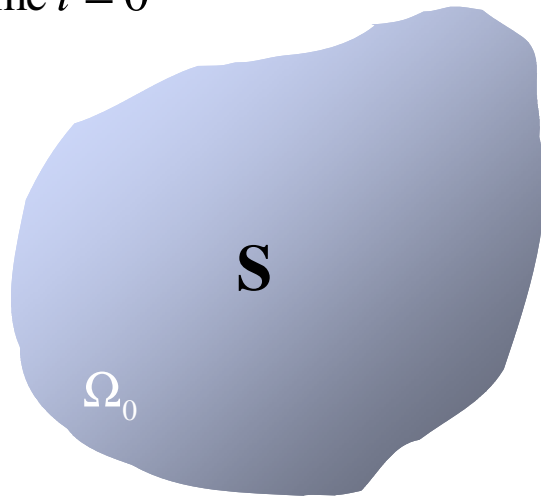
$$\boldsymbol{\tau} = \boldsymbol{\varphi}_* (\mathbf{S}) = \mathbf{F} \mathbf{S} \mathbf{F}^T$$

The *second Piola-Kirchhoff stress* tensor may be viewed as the **pull-back** of the (contravariant) *Kirchhoff stress* tensor, satisfying,

$$\mathbf{S} = \boldsymbol{\varphi}^* (\boldsymbol{\tau}) = \mathbf{F}^{-1} \boldsymbol{\tau} \mathbf{F}^{-T}$$

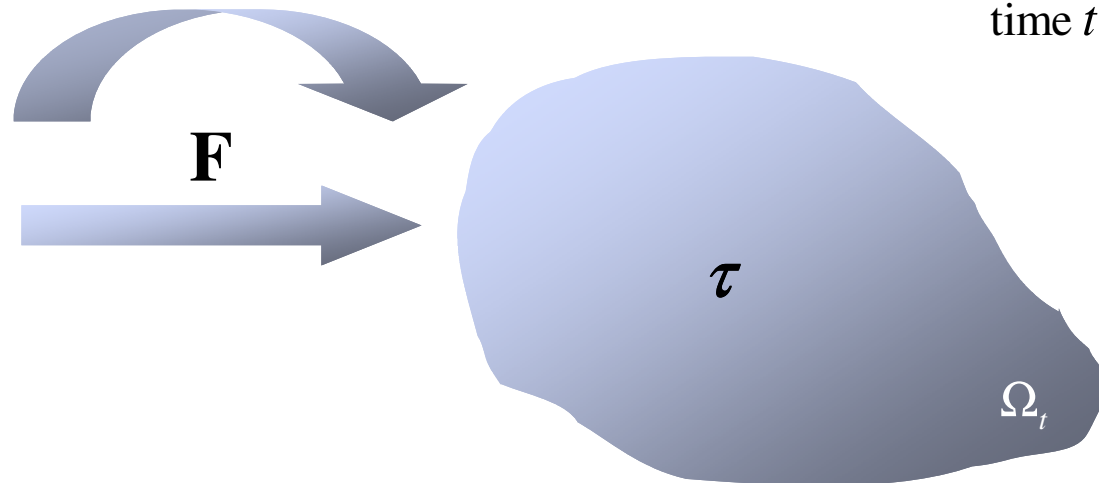
# Push-forward / Pull-back Maps

Reference or Material Configuration  
time  $t = 0$



$\varphi$

Current or Spatial Configuration  
time  $t$

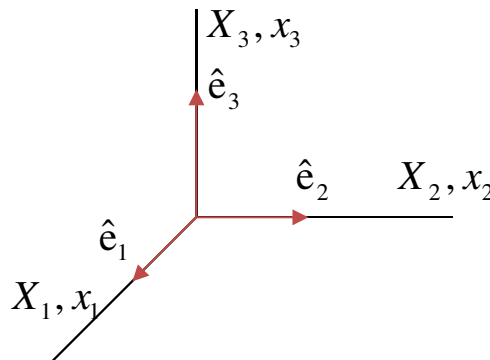


**Pull-back Maps**

$$\mathbf{S} = \mathbf{F}^{-1} \boldsymbol{\tau} \mathbf{F}^{-T}$$

**Push-forward Maps**

$$\boldsymbol{\tau} = \mathbf{F} \mathbf{S} \mathbf{F}^T$$



# Stress Tensors

## Cauchy Stress Tensor

$$\boldsymbol{\sigma} = J^{-1}\boldsymbol{\tau} = J^{-1}\mathbf{P}\mathbf{F}^T = J^{-1}\mathbf{F}\mathbf{S}\mathbf{F}^T$$

## Kirchhoff Stress Tensor

$$\boldsymbol{\tau} = J\boldsymbol{\sigma} = \mathbf{P}\mathbf{F}^T = \mathbf{F}\mathbf{S}\mathbf{F}^T$$

## First Piola-Kirchhoff Stress Tensor

$$\mathbf{P} = J\boldsymbol{\sigma}\mathbf{F}^{-T} = \boldsymbol{\tau}\mathbf{F}^{-T} = \mathbf{F}\mathbf{S}$$

## Second Piola-Kirchhoff Stress Tensor

$$\mathbf{S} = J\mathbf{F}^{-1}\boldsymbol{\sigma}\mathbf{F}^{-T} = \mathbf{F}^{-1}\boldsymbol{\tau}\mathbf{F}^{-T} = \mathbf{F}^{-1}\mathbf{P}$$